## A VARIATIONAL APPROACH TO THE EIGENFUNCTIONS OF THE ONE PARTICLE RELATIVISTIC HAMILTONIAN

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ABSTRACT. In this note we give a variational characterization of the eigenvalues and eigenvectors for the operator

$$H = H_0 + V = \sqrt{-c^2 \Delta + m^2 c^4} + V,$$

where  $H_0$  is the relativistic (free) Hamiltonian operator and V is a real valued potential. Our results hold when  $V(x) = -\frac{1}{|x|}$  and H describe a relativistic atom.

The characterization we give for the eigenvectors is useful in proving regularity and exponential decay of the solutions — properties which have been object of investigation by B. Simon with different techniques.

#### 1. INTRODUCTION AND MAIN RESULTS

In this note we give a variational characterization of the eigenvalues and eigenvectors (see Theorem 1) for the operator

$$H = H_0 + V = \sqrt{-c^2 \Delta + m^2 c^4} + V,$$

where  $H_0$  is the relativistic (free) Hamiltonian operator – which has been used to study models where relativistic effects became relevant – and V is a real valued potential. Our results hold when  $V(x) = -\frac{1}{|x|}$  and H describe a relativistic atom.

The characterization we give for the eigenvectors is useful in proving properties — such as regularity (see Theorem 3) and exponential decay of the solutions (see Theorem 2) — which have been object of investigation by B. Simon with different techniques in [16].

In order to describe our results, let us recall that to the operator  $H_0$  can be defined for all  $f \in H^1(\mathbb{R}^3)$  as the inverse Fourier transform of the  $L^2$  function  $\sqrt{c^2|p|^2 + m^2c^4} \hat{f}(p)$  (where  $\hat{f}$  denotes the Fourier transform of f). To  $H_0$  we can associate the following quadratic form

$$Q(f,g) = \int_{\mathbb{R}^3} \sqrt{c^2 |p|^2 + m^2 c^4} \, \hat{f}(p) \hat{g}(p) \, dp$$

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which can be extended to all functions  $f, g \in H^{1/2}(\mathbb{R}^3)$  where

$$H^{1/2}(\mathbb{R}^3) = \left\{ f \in L^2(\mathbb{R}^3) \mid \int_{\mathbb{R}^3} (1+|p|) |\hat{f}(p)|^2 \, dp < +\infty \right\}.$$

see for example [13] for more details.

On the potential V we assume

(h1)  $V \in L^3_w(\mathbb{R}^3) + L^\infty(\mathbb{R}^3), V \in L^\infty(\mathbb{R}^3 \setminus B_{R_0})$  for some  $R_0 > 0$  and (i)  $\lim_{R \to +\infty} ||V||_{L^\infty(|x|>R)} = 0;$ (ii)  $\lim_{R \to +\infty} \sup \operatorname{ess}_{|x|>R} V(x)|x|^2 = -\infty.$ 

(h2) V is  $H_0$  - form bounded with bound less than 1, i.e. there exists  $a \in (0,1)$  such that

$$|(\phi, V\phi)_{L^2}| \le a(\phi, H_0\phi)_{L^2}$$

for all 
$$\phi \in H^{1/2}(\mathbb{R}^3; \mathbb{C});$$

*Remark* 1. The above assumptions are similar to those used in the study of the characterization and computation of the eigenvalues for the Dirac-Coulomb Hamiltonian, to which our problem is related, see [8, 9] and references therein.

Remark 2. We recall that  $L^q_w(\mathbb{R}^N)$ , the weak  $L^q$  space, is the space of all measurable functions f such that

$$\sup_{\alpha>0} \alpha |\{ x \mid |f(x)| > \alpha \}|^{1/q} < +\infty,$$

where |E| denotes the Lebesgue measure of a measurable set  $E \subset \mathbb{R}^N$ . Note that  $f(x) = |x|^{-1}$  does not belong to any  $L^q$ -space but it belongs to  $L^3_w(\mathbb{R}^3)$ . (see e.g. [13] for more details).

Remark 3. The validity of (h2) when V is the Coulomb potential of a nucleus with Z protons

(1.1) 
$$V(x) = -\frac{Ze^2}{|x|} \quad (in \ cgs \ units)$$

follows from important inequalities. Let us recall them here.

**Hardy:** for all  $\psi \in H^1(\mathbb{R}^3)$ 

$$\left\| |x|^{-1}\psi \right\|_{{}_{L^2}} \le 2 \left\| \nabla \psi \right\|_{{}_{L^2}} \le \frac{2}{c\hbar} \left\| \sqrt{-c^2\hbar^2\Delta + m^2c^4}\psi \right\|_{{}_{L^2}}$$

**Kato, Herbst** [10]: for all  $\psi \in H^{1/2}(\mathbb{R}^3)$ 

$$\left(\psi, |x|^{-1}\psi\right)_{L^2} \le \frac{\pi}{2} \left(\psi, \sqrt{-\Delta}\psi\right)_{L^2} \le \frac{\pi}{2c\hbar} \left(\psi, \sqrt{-c^2\hbar^2\Delta + m^2c^4}\psi\right)_{L^2}$$

Note that (h2) is satisfied for the electrostatic potential provided 0 < Z < 68 by Hardy and provided 0 < Z < 87 by Kato.

Let us recall that the operator  $\sqrt{-c^2\Delta + m^2c^4}$ , exactly as for the fractional Laplacian, can be related, following [3], to a Dirichlet to Neumann operator (see also [2] and [5, 6, 7] for more closely related models).

To show this, we take a given function  $u \in \mathcal{S}(\mathbb{R}^3)$  with Fourier transform  $\hat{u}$  and let

$$v(x,y) = \frac{1}{(2\pi)^{3/2}} \int_{\mathbb{R}^3} e^{ip \cdot y} \hat{u}(p) e^{-\sqrt{c^2 |p|^2 + m^2 c^4} x} \, dp.$$

be the solution of the Dirichlet boundary problem

$$\begin{cases} -\partial_x^2 v - c^2 \Delta_y v + m^2 c^4 v = 0 & \text{in } \mathbb{R}^4_+ = \left\{ (x, y) \in \mathbb{R} \times \mathbb{R}^3 \mid x > 0 \right\} \\ v(0, y) = u(y) & \text{for } y \in \mathbb{R}^3 = \partial \mathbb{R}^4_+. \end{cases}$$

Setting

$$\mathcal{T}u(y) = \frac{\partial v}{\partial \nu}(0, y) = -\frac{\partial v}{\partial x}(0, y)$$

we have that

$$\mathcal{T}u(y) = -\frac{\partial v}{\partial x}(0, y) = \frac{1}{(2\pi)^{3/2}} \int_{\mathbb{R}^3} e^{ip \cdot y} \sqrt{c^2 |p|^2 + m^2 c^4} \,\hat{u}(p) \, dp$$

namely  $\mathcal{T} = \sqrt{-c^2 \Delta_y + m^2 c^4} = H_0$  on the dense domain  $\mathcal{S}(\mathbb{R}^3)$ . We consider the functional  $\mathcal{I}(\phi)$  defined on  $H^1(\mathbb{R}^4_+, \mathbb{C})$ 

(1.2) 
$$\mathcal{I}(\phi) = \iint_{\mathbb{R}^4_+} (|\partial_x \phi|^2 + c^2 |\nabla_y \phi|^2 + m^2 c^4 |\phi|^2) \, dx \, dy + \int_{\mathbb{R}^3} (\phi_{tr}, V \phi_{tr}) \, dy$$

where  $\phi_{tr} \in H^{1/2}$  denotes the trace of  $\phi \in H^1$  on  $\partial \mathbb{R}^4_+ = \mathbb{R}^3$ .

We have the following existence and characterization results for the eigenvalues and eigenfunctions, where we always assume m > 0.

**Theorem 1.** Let m > 0 and (h1)-(h2) hold. Then there exist  $\lambda_1 \leq \lambda_2 \leq \ldots \leq \lambda_k \leq \ldots$  and  $\phi_1, \phi_2, \ldots, \phi_k, \ldots \in H^1(\mathbb{R}^4_+, \mathbb{C})$  such that, for all  $k \in \mathbb{N}$ 

$$\lambda_k = \mathcal{I}(\phi_k) = \inf_{X_k} \mathcal{I}(\phi)$$

where

$$X_1 = \left\{ \phi \in H^1 \mid |\phi_{tr}|_{L^2} = 1 \right\}.$$

and, for  $1 < k \in \mathbb{N}$ 

$$X_k = \left\{ \phi \in H^1 \mid |\phi_{tr}|_{L^2} = 1, \ (\phi_{tr}, (\phi_i)_{tr})_{L^2} = 0, \ i = 1, \dots, k-1 \right\}.$$

Moreover  $\{\lambda_k\}_{k\geq 1} \in \sigma_{disc}(H_0+V)$  and

$$0 < \lambda_1 \leq \ldots \leq \lambda_k \leq \lambda_{k+1} \rightarrow \inf\{\sigma_{ess}(H_0 + V)\} = mc^2 \quad for \ k \rightarrow +\infty.$$

The functions  $\varphi_k = (\phi_k)_{tr} \in H^{1/2}(\mathbb{R}^3, \mathbb{C})$  are the eigenfunctions of the operator  $H_0 + V$ , and the functions  $\phi_k \in H^1(\mathbb{R}^4_+, \mathbb{C})$  are weak solution of the Neumann problem

$$(\mathcal{E}_k) \qquad \begin{cases} -\partial_x^2 \phi_k - c^2 \Delta_y \phi_k + m^2 c^4 \phi_k = 0 & \text{ in } \mathbb{R}^4_+ \\ \frac{\partial \phi_k}{\partial \nu} + V \varphi_k = \lambda_k \varphi_k & \text{ on } \partial \mathbb{R}^4_+ = \mathbb{R}^3 \end{cases}$$

The following Theorems give some properties of the eigenfunctions: regularity and exponential decay.

**Theorem 2** (exponential decay). Let m > 0 and  $(h_1)$ - $(h_2)$  hold. Let  $\phi_k \in H^1(\mathbb{R}^4_+, \mathbb{C})$  (and  $\varphi_k = (\phi_k)_{tr}$ ) be the functions given by the Theorem 1.

Then for all  $0 \leq \beta < \sqrt{m^2 c^4 - \lambda_k^2}$  there exists R > 0 such that  $e^{\frac{\beta}{c}|y|} \varphi_k \in L^2(\mathbb{R}^3 \setminus B_R)$ .

*Remark* 4. Several authors have investigated the asymptotic behaviour of eigenfunctions. Let us recall here the classical book by Agmon [1] and [14, 15, 4].

**Theorem 3** (regularity). Let  $\phi_k \in H^1(\mathbb{R}^4_+, \mathbb{C})$  (and  $\varphi_k = (\phi_k)_{tr}$ ) be the functions given by the Theorem 1 and  $R_0$  be given by (h1).

Then we have

- (i)  $\phi_k \in W^{1,q}([0,r) \times (\mathbb{R}^3 \setminus B_{R_0}))$  for any  $q \in [2,\infty]$ , r > 0; (ii)  $\phi_k \in C^{0,\alpha}([0,+\infty) \times (\mathbb{R}^3 \setminus B_{R_0}))$  for any  $\alpha \in [0,1]$  and  $\varphi_k \in C^{0,\alpha}(\mathbb{R}^3 \setminus B_{R_0})$ ; (iii) if in addition  $V \in L^3_{loc}(\mathcal{U})$  for some  $\mathcal{U} \subset \mathbb{R}^3$  then for every  $\mathcal{V} \subset \subset \mathcal{U}$  (i.e. such that its closure is compact in  $\mathcal{U}$   $\phi_k \in W^{1,p}([0,r) \times \mathcal{V})$  for any  $p \in [2,\infty)$  and r > 0 and  $\varphi_k \in C^{0,\alpha}(\mathcal{V})$  for any  $\alpha \in [0,1)$ .

### 2. Proof of Theorem 1

We divide the proof of Theorem 1 in several steps.

2.1. Notation and preliminary results. With  $||u||_p$  we will denote the norm of  $u \in$  $L^p(\mathbb{R}^4_+)$  and with  $|v|_p$  the norm of  $v \in L^p(\mathbb{R}^3)$ .

We introduce the following (equivalent) norm in  $H^1(\mathbb{R}^4_+,\mathbb{C})$ 

$$\|\phi\|_{H^1}^2 = \iint_{\mathbb{R}^4_+} (|\partial_x \phi|^2 + c^2 |\nabla_y \phi|^2 + m^2 c^4 |\phi|^2) \, dx \, dy.$$

and the following norm in the weak  $L^q$ -space:

$$|f|_{L^{q}_{w}} = \sup \left\{ |A|^{-1/r} \int_{A} |f(x)| \, dx \mid A \subset \mathbb{R}^{3}, \text{measurable}, \ 0 < |A| < +\infty \right\}$$

where 1/q + 1/r = 1. For the weak  $L^q$  spaces the following generalization of the weak Young inequality holds:

**Proposition 1** (see [11, thm. 2.10]). Let  $f \in L^q_w(\mathbb{R}^N)$ ,  $g \in L^{q'}_w(\mathbb{R}^N)$  and  $h \in L^p(\mathbb{R}^N)$  with  $\frac{1}{q} + \frac{1}{q'} = 1$  and 1 . Then

(2.1) 
$$||f(g*h)||_{p} \le C||f||_{q,w} ||g||_{q',w} ||h||_{p}.$$

From this we deduce the following result:

Lemma 1. Let  $V \in L^3_w(\mathbb{R}^3)$  and  $f \in H^{1/2}(\mathbb{R}^3)$ . Then

(2.2) 
$$||V|^{1/2}f|_{L^2} \le C|V|_{L^3_w}^{1/2}|f|_{H^{1/2}}.$$

*Proof.* Follows from [12, (42)] that the Green function  $G^{\mu}_{\alpha}$  of  $(-\Delta + \mu^2)^{\alpha/2}$  belongs to  $L_w^{3/(3-\alpha)}(\mathbb{R}^3)$  if  $\mu \ge 0$  and  $0 < \alpha < 3$ .

Then, given  $f \in H^{1/2}(\mathbb{R}^3)$ , let  $h = (-\Delta + \mu^2)^{1/4} f \in L^2(\mathbb{R}^3)$ ,  $f = G^{\mu}_{1/2} * h$ . From the weak Young's inequality above (2.1), we deduce

$$\begin{split} ||V|^{1/2}f|_{L^2} &= ||V|^{1/2} (G_{1/2}^{\mu} * h)|_{L^2} \le C ||V|^{1/2}|_{L_w^6} |G_{1/2}^{\mu}|_{L_w^{6/5}} |h|_{L^2} \\ &\le C |V|_{L_w^3} |(-\Delta + \mu^2)^{1/4} f|_{L^2} \le C |V|_{L_w^3}^{1/2} |f|_{H^{1/2}}. \end{split}$$

We also recall that for all  $v \in C_0^{\infty}(\mathbb{R}^4)$ 

$$\int_{\mathbb{R}^3} |v(0,y)|^2 dy = \int_{\mathbb{R}^3} dy \int_{+\infty}^0 \partial_x |v|^2 dx \le 2 \|v\|_{L^2(\mathbb{R}^4_+)} \|\partial_x v\|_{L^2(\mathbb{R}^4_+)}$$

and by density we get for all  $\phi \in H^1$  and any  $\alpha > 0$ 

(2.3) 
$$\alpha \int_{\mathbb{R}^3} |\phi_{\rm tr}|^2 \, dy \leq \iint_{\mathbb{R}^4_+} (|\partial_x \phi|^2 + \alpha^2 |\phi|^2) \, dx \, dy.$$

This implies in particular that the quadratic form (kinetic energy)

(2.4) 
$$\mathcal{T}(\phi) = \iint_{\mathbb{R}^4_+} (|\partial_x \phi|^2 + m^2 c^4 |\phi|^2) \, dx \, dy - m c^2 |\phi_{tr}|_{L^2}^2 \ge 0$$

is positive definite.

We introduce the differential  $d\mathcal{I}(\phi) \colon H^1 \to \mathbb{R}$  of the functional  $\mathcal{I}$ 

$$d\mathcal{I}(\phi)[h] = 2\operatorname{Re} \iint_{\mathbb{R}^4_+} \left( (\partial_x \phi, \partial_x h) + c^2 \left( \nabla_y \phi, \nabla_y h \right) + m^2 c^4 \left( \phi, h \right) \right) \, dx \, dy \\ + 2\operatorname{Re}(\phi_{tr}, Vh_{tr})_{L^2}$$

The following property can be easily verified.

**Lemma 2.** For  $w \in H^1(\mathbb{R}^4_+)$ , let  $u = w_{tr} \in H^{1/2}(\mathbb{R}^3)$  be the trace of w,  $\hat{u} = \mathcal{F}(u)$  and

$$v(x,y) = \mathcal{F}_y^{-1} \big[ \hat{u}(p) e^{-\sqrt{c^2 |p|^2 + m^2 c^4 x}} \big].$$

Then 
$$v \in H^{1}(\mathbb{R}^{4}_{+})$$
,  $||v||_{H^{1}(\mathbb{R}^{4})} = ||u||_{H^{1/2}(\mathbb{R}^{3})}$ , and  

$$\int_{\mathbb{R}^{3}} \sqrt{c^{2}|p|^{2} + m^{2}c^{4}} |\hat{u}|^{2} dp = \iint_{\mathbb{R}^{4}_{+}} (|\partial_{x}v|^{2} + c^{2}|\nabla_{y}v|^{2} + m^{2}c^{4}|v|^{2}) dx dy$$

$$\leq \iint_{\mathbb{R}^{4}_{+}} (|\partial_{x}w|^{2} + c^{2}|\nabla_{y}w|^{2} + m^{2}c^{4}|w|^{2}) dx dy.$$

In other words

(2.5) 
$$\|w_{tr}\|_{H^{1/2}}^2 = (w_{tr}, H_0 w_{tr})_{L^2} \le \|w\|_{H^1}^2$$
 for every  $w \in H^1(\mathbb{R}^4_+)$ 

2.2. Existence of the ground state. We consider the following minimization problem :  $(\mathcal{P}_1)$   $\lambda_1 = \inf_{\phi \in S} \mathcal{I}(\phi).$ 

where  $S = \left\{ \phi \in H^1 \ \Big| \ |\phi_{tr}|_{L^2}^2 = 1 \right\}.$ 

Lemma 3. The following holds:

(i) *I*(φ) is bounded by below and coercive on H<sup>1</sup>,
(ii) 0 < λ<sub>1</sub> < mc<sup>2</sup>.

*Proof.* (i) Let  $\phi \in H^1$ ,  $\varphi = \phi_{tr}$ . From (**h2**) and (2.5), there exists  $a \in (0, 1)$  such that

$$\begin{split} (\varphi, V\varphi)_{L^2} &\geq -a(\varphi, \sqrt{-c^2\Delta + m^2c^4}\,\varphi)_{L^2} \\ &\geq -a \iint_{\mathbb{R}^4_+} (|\partial_x \phi|^2 + c^2 |\nabla_y \phi|^2 + m^2c^4 |\phi|^2) \, dx \, dy \end{split}$$

Therefore, we may conclude that there exists  $\delta > 0$  such that  $\mathcal{I}(\phi) \geq \delta \|\phi\|_{\mu^1}^2$ .

(ii) From (i) immediately follows that  $\lambda_1 > 0$ . Now take  $\phi(x, y) = e^{-mc^2 x} \varphi(y)$ , with  $\varphi \in C_0^{\infty}(\mathbb{R}^3, \mathbb{C})$ , and  $|\varphi|_{L^2} = 1$ , we have

$$\mathcal{I}(\phi) - mc^2 = \frac{1}{2m} \int_{\mathbb{R}^3} |\nabla \varphi|^2 + \int_{\mathbb{R}^3} V |\varphi|^2 \, dy = \mathcal{E}(\varphi)$$

Take, now  $\varphi_{\eta}(y) = \eta^{3/2} \varphi(\eta y)$ , we have  $|\varphi_{\eta}|_{L^2} = 1$ , for any  $\eta > 0$  and setting  $\phi_{\eta}(x, y) = e^{-mc^2 x} \varphi_{\eta}(y)$ 

$$\begin{split} \lambda_1 - mc^2 &\leq \inf_{\eta > 0} \mathcal{I}(\phi_\eta) - mc^2 = \inf_{\eta > 0} \mathcal{E}(\varphi_\eta) = \\ &= \inf_{\eta > 0} \eta^2 \frac{1}{2m} \int_{\mathbb{R}^3} |\nabla \varphi|^2 \, dy + \int_{\mathbb{R}^3} V(\eta^{-1}y) |\varphi|^2 \, dy \end{split}$$

By (h1), for any K > 0 there exists R > 0 such that for any |y| > R we have  $V(y) \le -K/|y|^2$  a.e.. Hence

$$\begin{split} \left(\varphi, V(\eta^{-1}y)\varphi\right)_{L^2} &= \int_{\{\eta^{-1}|y| \le R\}} V(\eta^{-1}y)|\varphi|^2 + \int_{\{\eta^{-1}|y| > R\}} V(\eta^{-1}y)|\varphi|^2 \\ &\leq \eta^3 \sup_{|y| \le \eta R} |\varphi(y)|^2 \int_{\{|y| \le R\}} |V(y)| - K\eta^2 \int_{\{|y| > \eta R\}} \frac{1}{|y|^2} |\varphi|^2 \\ &\leq C(\eta^3 - K\eta^2) \end{split}$$

where the constant C > 0 depends on  $\varphi$  and R, and K > 0 is arbitrarily large.

We immediately conclude that for any given  $\varphi\in C_0^\infty(\mathbb{R}^3;\mathbb{C})$ 

$$\limsup_{\eta \to 0^+} \frac{1}{\eta^2} (\varphi, V(\eta^{-1}y)\varphi)_{L^2} = -\infty$$
  
$$0 < 0.$$

which implies that  $\lambda_1 - mc^2 < 0$ .

Letting  $\mathcal{G}(\phi) = |\phi_{tr}|_{L^2}^2$  we have that  $S = \{\phi \in H^1 \mid \mathcal{G}(\phi) = 1\}$  and the tangent space at S at the point  $\phi \in S$  is the set

$$T_{\phi}S = \left\{ h \in H^1 \mid d\mathcal{G}(\phi)[h] = 2 \operatorname{Re}(\phi_{tr}, h_{tr})_{L^2} = 0 \right\}$$

and that  $\nabla_S \mathcal{I}(\phi)$ , the projection of the gradient on the tangent space  $T_{\phi}S$  to S at the point  $\phi$  is given by

$$\nabla_S \mathcal{I}(\phi) = \nabla \mathcal{I}(\phi) - \mu(\phi) \nabla \mathcal{G}(\phi)$$

where  $\nabla \mathcal{I}(\phi) \in H^1$  is such that

$$\left(\nabla \mathcal{I}(\phi), h\right)_{H^1} = d\mathcal{I}(\phi)[h] = 2\operatorname{Re}(\phi, h)_{H^1} + 2\operatorname{Re}(\phi_{tr}, Vh_{tr})_{L^2} \quad \text{for all } h \in H^1,$$

 $\nabla \mathcal{G}(\phi) \in H^1$  is such that

$$(\nabla \mathcal{G}(\phi), h)_{H^1} = d\mathcal{G}(\phi)[h] = 2 \operatorname{Re}(\phi_{tr}, h_{tr})_{L^2}$$
 for all  $h \in H^1$ ,  
and  $\mu(\phi) \in \mathbb{R}$  is such that  $\nabla_S \mathcal{I}(\phi) \in T_{\phi}S$ . Then

$$0 = \left(\nabla \mathcal{G}(\phi), \nabla_S \mathcal{I}(\phi)\right)_{H^1} = \left(\nabla \mathcal{G}(\phi), \nabla \mathcal{I}(\phi)\right)_{H^1} - \mu(\phi) \|\nabla \mathcal{G}(\phi)\|_{H^1}^2$$

and

$$\mu(\phi) = \frac{\left(\nabla \mathcal{G}(\phi), \nabla_S \mathcal{I}(\phi)\right)_{H^1}}{\|\nabla \mathcal{G}(\phi)\|_{H^1}^2}$$

From

$$(\nabla_{S} \mathcal{I}(\phi), \phi)_{H^{1}} = (\nabla \mathcal{I}(\phi), \phi)_{H^{1}} - \mu(\phi) (\nabla \mathcal{G}(\phi), \phi)_{H^{1}}$$
  
=  $2\mathcal{I}(\phi) - 2\mu(\phi)\mathcal{G}(\phi) = 2\mathcal{I}(\phi) - 2\mu(\phi)$ 

we also deduce that

(2.6) 
$$\mu(\phi) = \mathcal{I}(\phi) - \frac{1}{2} (\nabla_S \mathcal{I}(\phi), \phi)_{\mu}$$

We now recall the following well known result

**Lemma 4.** There exists a Palais-Smale minimizing sequence  $\phi_n$  for  $\mathcal{I}$  on the set  $S = \{\phi \mid$  $|\phi_{tr}|_{L^2}^2 = 1$ , that is a sequence such that, denoting  $\varphi_n = (\phi_n)_{tr}$ ,

$$\mathcal{I}(\phi_n) \to \lambda_1, \qquad \nabla_S \mathcal{I}(\phi_n) \to 0, \qquad |\varphi_n|_{L^2}^2 = 1$$

*Proof.* Assuming that the result does not hold, one deduces that there exist  $\epsilon > 0, \delta > 0$ such that  $\|\nabla_S \mathcal{I}(\phi)\| \ge \delta > 0$  for all  $\phi \in S$  such that  $\lambda_1 - \epsilon < \mathcal{I}(\phi) < \lambda_1 + \epsilon$ . Then one can build a gradient flow  $\eta' = \nabla_S \mathcal{I}(\eta)$ , which leaves S invariant and pushes  $\{\mathcal{I} < \lambda_1 + \epsilon\} \cap S$ into  $\{\mathcal{I} < \lambda_1 - \epsilon\} \cap S$ , a contradiction.

The Lemma also follows from Ekeland's variational principle.

**Lemma 5.** Let  $\phi_n$  be a Palais Smale sequence at some level  $\lambda \geq 0$  for  $\mathcal{I}$  on S. Let  $\varphi_n = (\phi_n)_{tr}.$ If  $\varphi_n \rightharpoonup 0$  in  $H^{1/2}$  then

$$(\varphi_n, V\varphi_n)_{L^2} \to 0.$$

*Proof.* Since  $\mathcal{I}$  is coercive,  $\phi_n$  is bounded  $H^1$ ,  $\varphi_n$  is bounded in  $H^{1/2}$  and, by Sobolev embedding, relatively compact in  $L_{\text{loc}}^p$  for  $p \in [2,3)$ . From (2.6) follows that also  $\mu_n$  is bounded.

By (h1)  $V \in L^{\infty}(\mathbb{R}^3 \setminus B_{R_0})$  and for any  $\varepsilon > 0$ , the set  $A_{\varepsilon} = \{ y \in \mathbb{R}^3 \setminus B_{R_0} \mid |V(y)| \ge \varepsilon \}$ is bounded.

Take a radial function  $\chi \in C_0^{\infty}(\mathbb{R}^3)$ , with values in [0,1] such that  $\chi(y) = 1$  for  $y \in B_1$ and  $\chi(y) = 0$  for  $y \in \mathbb{R}^3 \setminus B_2$  and let  $\chi_R(y) = \chi(R^{-1}y)$ .

Taking  $R > R_0$  in such a way that  $A_{\varepsilon} \subset B_R$  we have

$$\left|\left(\varphi_n, (1-\chi_R^2)V\varphi_n\right)_{L^2}\right| \le \epsilon |\varphi_n|_{L^2}^2 \le \varepsilon.$$

We have, by assumption,  $\mathcal{I}(\phi_n) \to \lambda$ ,  $\mathcal{G}(\phi_n) = |\varphi_n|_{r^2}^2 = 1$  and

$$\|\nabla_S \mathcal{I}(\phi_n)\| = \|d\mathcal{I}(\phi_n) - \mu_n d\mathcal{G}(\phi_n)\| \to 0.$$

where  $\mu_n = \mu(\phi_n)$  and also, by (2.6)

(2.7) 
$$\mu_n = \mathcal{I}(\phi_n) - \frac{1}{2} (\nabla_S \mathcal{I}(\phi_n), \phi_n)_{H^1} \to \lambda$$

Since  $\chi$  depends only on y we have that

(2.8) 
$$d\mathcal{I}(\phi_n)[\chi_R^2\phi_n] = d\mathcal{I}(\chi_R\phi_n)[\chi_R\phi_n] - 2c^2 \|\varphi_n\nabla_y\chi_R\|_{L^2}^2$$

and since  $C \|\phi_n\|_{H^1} \ge \|\chi_R^2 \phi_n\|_{H^1}$  we have that

 $\Box$ 

$$\begin{split} o_n(1) &= C \|\nabla \mathcal{I}(\phi_n) - \mu_n \nabla \mathcal{G}(\phi_n)\| \|\phi_n\|_{H^1} \ge |(\nabla \mathcal{I}(\phi_n) - \mu_n \nabla \mathcal{G}(\phi_n), \chi_R^2 \phi_n)_{H^1}| \\ &\ge |d\mathcal{I}(\phi_n)[\chi_R^2 \phi_n]| - |\mu_n 2 \operatorname{Re}(\varphi_n, \chi_R^2 \varphi_n)_{L^2}| \\ &\ge 2\mathcal{I}(\chi_R \phi_n) - 2c^2 \|\phi_n \nabla_y \chi_R\|_{L^2}^2 - 2|\mu_n| |\chi_R \varphi_n|_{L^2}^2 \end{split}$$

Now, by Sobolev compact embedding, for any given R > 0,

 $|\chi_{\scriptscriptstyle R} \varphi_n|_{_{L^2}} \to 0 \qquad \text{as} \quad n \to +\infty.$ 

Moreover,

$$\|\phi_n \nabla \chi_R\|_{L^2}^2 \le C \sup_{y \in \mathbb{R}^3} |\nabla \chi_R|^2 \le \frac{C}{R^2}$$

Since by Lemma 3-(i) we have

$$\mathcal{I}(\chi_{R}\phi_{n}) \geq \delta \|\chi_{R}\phi_{n}\|_{H^{1}}^{2}$$

we may conclude (recalling that  $\mu_n$  is bounded) that

$$\|\chi_{R}\phi_{n}\|_{H^{1}}^{2} \leq \epsilon_{n} + \frac{C}{R}.$$

and hence by (h2) and (2.5) we get

$$\left|\left(\chi_{R}\varphi_{n}, V\chi_{R}\varphi_{n}\right)_{L^{2}}\right| \leq a\left|\left(\chi_{R}\varphi_{n}, H_{0}\chi_{R}\varphi_{n}\right)_{L^{2}}\right| \leq a\left\|\chi_{R}\phi_{n}\right\|_{H^{1}}^{2} \leq \epsilon_{n} + \frac{C}{R}$$

for some  $\epsilon_n \to 0$  and R arbitrarily large.

Now we may conclude the existence of a minimizer for  $\mathcal{P}_{_1}.$  We have the following Proposition:

**Proposition 2.** Let  $\phi_n$  be a minimizing Palais Smale sequence at level  $\lambda_1 > 0$  for  $\mathcal{I}$  with  $|(\phi_n)_{tr}|_{L^2} = 1$  (as in Lemma 5).

Then  $\phi_n \rightarrow \phi \not\equiv 0$  in  $H^1$  and  $\tilde{\phi} = |\phi_{tr}|_{L^2}^{-1} \phi$  is a minimizer for  $\mathcal{I}$  on S, that is

$$\mathcal{I}(\phi) = \lambda_1, \qquad |\phi_{tr}|_{L^2} = 1.$$

Moreover  $\tilde{\phi}$  (and hence also  $\phi$ ) is a weak solution of the Neumann problem  $(\mathcal{E}_{_1})$ .

*Proof.* Since  $\mathcal{I}$  is coercive,  $\phi_n$  is bounded (and weakly convergent) in  $H^1$ ,  $\varphi_n = (\phi_n)_{tr}$  is bounded (and weakly convergent) in  $H^{1/2}$ .

If by contradiction  $\varphi_n \rightharpoonup \varphi \equiv 0$ , then by Lemma 5 we have

$$(\varphi_n, V\varphi_n)_{L^2} \to 0$$

Now, by (2.4) we get

$$\mathcal{I}(\phi_n) - mc^2 |\varphi_n|_{L^2}^2 \ge (\varphi_n, V\varphi_n)_{L^2} \to 0.$$

On the other hand, by Lemma 3-(ii)

$$\mathcal{I}(\phi_n) - mc^2 |\varphi_n|_{L^2}^2 = \mathcal{I}(\phi_n) - mc^2 \to \lambda_1 - mc^2 < 0$$

a contradiction, that is  $\varphi_n \rightharpoonup \varphi \not\equiv 0$ .

It follows from (2.6) that

$$\mu_n = \mathcal{I}(\phi_n) - \frac{1}{2} (\nabla_S \mathcal{I}(\phi_n), \phi_n)_{H^1} \to \lambda_1$$

and hence, by weak convergence, we have

$$d\mathcal{I}(\phi_n)[h] - \mu_n d\mathcal{G}(\phi_n)[h] \to d\mathcal{I}(\phi)[h] - \lambda_1 d\mathcal{G}(\phi)[h] = 0 \quad \forall h \in H^1$$

hence in particular

$$0 = d\mathcal{I}(\phi)[\phi] - \lambda_1 d\mathcal{G}(\phi)[\phi] = 2\mathcal{I}(\phi) - 2\lambda_1 \mathcal{G}(\phi)$$

and we may conclude that  $\tilde{\phi} = \mathcal{G}(\phi)^{-1/2} \phi$  is a minimizer for  $\mathcal{I}$  on S, namely

$$\lambda_1 = \frac{\mathcal{I}(\phi)}{\mathcal{G}(\phi)} = \mathcal{I}(\mathcal{G}(\phi)^{-1/2}\phi) = \mathcal{I}(\tilde{\phi})$$
$$\mathcal{G}(\tilde{\phi}) = \mathcal{G}(\mathcal{G}(\phi)^{-1/2}\phi) = 1$$

Now, we look for the existence of higher eigenvalues and corresponding eigenfunctions. We proceed by induction.

Let  $\lambda_1$  be defined by  $(\mathcal{P}_1)$  and  $\phi_1$  be the corresponding minimizer given by Proposition 2.

Assume we have defined, for j = 1, ..., k - 1,  $\lambda_1 \leq \cdots \leq \lambda_j \leq \ldots \leq \lambda_{k-1} < mc^2$  and  $\phi_j \in H^1$ ,  $\varphi_j = (\phi_j)_{tr} \in H^{1/2}$  such that

$$(\varphi_i, \varphi_j)_{L^2} = \delta_{ij}, \qquad i, j = 1, \dots, k-1,$$

and

$$(\mathcal{P}_j)$$
  $\lambda_j = \mathcal{I}(\phi_j) = \inf_{\phi \in X_j} \mathcal{I}(\phi)$   $j = 1, \dots, k-1$ 

where,

$$X_{j} = \left\{ \phi \in H^{1} \mid \mathcal{G}(\phi) = |\phi_{tr}|_{L^{2}}^{2} = 1, \ (\phi_{tr}, \varphi_{i})_{L^{2}} = 0 \quad \text{for} \quad i = 1, \dots, j-1 \right\}.$$

We define

$$(\mathcal{P}_k) \qquad \qquad \lambda_k = \inf_{\phi \in X_k} \mathcal{I}(\phi)$$

Remark 5. Setting  $\mathcal{G}_j(\phi) = (\varphi_j, \phi_{tr})_{L^2}$ , for  $j \ge 1$ , we have that the linear functionals  $\mathcal{G}_j$  are bounded on  $H^1$  and for any  $\phi, h \in H^1$ 

$$d\mathcal{G}_{j}(\phi)[h] = \left(\nabla \mathcal{G}_{j}(\phi), h\right)_{H^{1}} = \left(\varphi_{j}, h_{tr}\right)_{L^{2}} = \mathcal{G}_{j}(h) \qquad j = 1, \dots k - 1.$$
  
Then  $X_{k} = \left\{\phi \in H^{1} \mid \mathcal{G}(\phi) = 1, \ \mathcal{G}_{j}(\phi) = 0, \ j = 1, \dots, k - 1\right\},$ 

$$T_{\phi}X_{k} = \left\{ h \in H^{1} \mid (\nabla \mathcal{G}(\phi), h)_{H^{1}} = 0, \ \mathcal{G}_{j}(h) = 0, \ j = 1, \dots, k-1 \right\}$$

and the constrained gradient (i.e. the projection of the gradient of  $\mathcal{I}$  on the tangent space  $T_{\phi}X_k$ ) is given by

$$\nabla_{X_k} \mathcal{I}(\phi) = \nabla \mathcal{I}(\phi) - \mu_0(\phi) \nabla \mathcal{G}(\phi) - \sum_{j=1}^{k-1} \mu_j(\phi) \nabla \mathcal{G}_j(\phi).$$

Taking  $\phi \in X_k$  we have that

$$(\nabla_{X_k} \mathcal{I}(\phi), \phi)_{H^1} = (\nabla \mathcal{I}(\phi), \phi)_{H^1} - \mu_0(\phi) (\nabla \mathcal{G}(\phi), \phi)_{H^1} - \sum_{j=1}^{k-1} \mu_j(\phi) (\nabla \mathcal{G}_j(\phi), \phi)_{H^1}$$
  
=  $2\mathcal{I}(\phi) - 2\mu_0(\phi)\mathcal{G}(\phi) - \sum_{j=1}^{k-1} \mu_j(\phi)\mathcal{G}_j(\phi) = 2\mathcal{I}(\phi) - 2\mu_0(\phi)$ 

and we deduce that

(2.9) 
$$\mu_0(\phi) = \mathcal{I}(\phi) - \frac{1}{2} (\nabla_{X_k} \mathcal{I}(\phi), \phi)_{\mu_1}$$

Taking again  $\phi \in X_k$  and  $i = 1, \ldots, k - 1$ , from

$$(\nabla_{X_{k}}\mathcal{I}(\phi),\phi_{i})_{H^{1}} = (\nabla\mathcal{I}(\phi),\phi_{i})_{H^{1}} - \mu_{0}(\phi)(\nabla\mathcal{G}(\phi),\phi_{i})_{H^{1}} - \sum_{j=1}^{k-1} \mu_{j}(\phi)(\nabla\mathcal{G}_{j}(\phi),\phi_{i})_{H^{1}}$$
$$= d\mathcal{I}(\phi)[\phi_{i}] - \mu_{0}(\phi)2\operatorname{Re}(\phi_{tr},\varphi_{i})_{L^{2}} - \sum_{j=1}^{k-1} \mu_{j}(\phi)(\varphi_{j},\varphi_{i})_{L^{2}}$$
$$= d\mathcal{I}(\phi)[\phi_{i}] - \mu_{i}(\phi)$$

we have that

(2.10) 
$$\mu_i(\phi) = d\mathcal{I}(\phi)[\phi_i] - \left(\nabla_{X_k} \mathcal{I}(\phi), \phi_i\right)_{H^2}$$

We say that  $\phi_n \in X_k$  is a (constrained) Palais Smale sequence for  $\mathcal{I}$  on  $X_k$  at level  $\lambda_k$  if  $\phi_n \in X_k$ ,

$$\mathcal{I}(\phi_n) \to \lambda_k \text{ and } \|\nabla_{x_k} \mathcal{I}(\phi_n)\| \to 0.$$

The proof of existence of a minimizer for  $(\mathcal{P}_k)$  proceeds as the proof of the existence of the ground state  $\phi_1$ . The key points are the following two Lemmas.

# Lemma 6. $\lambda_1 \leq \lambda_k < mc^2$ .

Proof. Let us consider any k-dimensional linear subspace  $G_k \subset C_0^{\infty}(\mathbb{R}^3; \mathbb{C})$ . For  $\varphi \in G_k \cap S$  and  $\eta > 0$  we let  $\varphi_{\eta}(y) = \eta^{3/2} \varphi(\eta y) \in S$  and

$$F_k^\eta = \left\{ \, \phi_\eta \in H^1 \ \Big| \ \phi_\eta(x,y) = \mathrm{e}^{-mc^2 x} \, \varphi_\eta(y), \quad \varphi \in G_k \cap S \, \right\}.$$

Then, for any  $\phi_{\eta} \in F_k^{\eta}$ 

$$\begin{split} \mathcal{I}(\phi_{\eta}) - mc^2 &= \frac{1}{2m} \int_{\mathbb{R}^3} |\nabla \varphi_{\eta}|^2 + \int_{\mathbb{R}^3} (\varphi_{\eta}, V\varphi_{\eta}) \\ &= \frac{\eta^2}{2m} \int_{\mathbb{R}^3} |\nabla \varphi|^2 + \int_{\mathbb{R}^3} V(\eta^{-1}y) |\varphi|^2 \end{split}$$

Arguing as in Lemma 3-(ii) and by compactness of the set  $G_k \cap S$ , there exists  $\bar{\eta} > 0$  such that for any  $\phi_{\bar{\eta}} \in F_k^{\bar{\eta}}$ , we have

$$\bar{\eta}^2 \frac{1}{2m} \int_{\mathbb{R}^3} |\nabla \varphi|^2 + \int_{\mathbb{R}^3} V(y/\bar{\eta}) |\varphi|^2 < 0$$

Since  $X_k \cap F_k^{\bar{\eta}} \neq \emptyset$ , we have  $\lambda_k \leq \sup_{F_k^{\bar{\eta}}} \mathcal{I}(\phi_{\bar{\eta}}) < mc^2$ .

**Lemma 7.** Let  $\zeta_n \in X_k$  be a (constrained) Palais Smale sequence at level  $\lambda_k$  for  $\mathcal{I}$  on  $X_k$ , with gradient

$$\nabla_{X_k} \mathcal{I}(\zeta_n) = \nabla \mathcal{I}(\zeta_n) - \mu_0(\zeta_n) \nabla \mathcal{G}(\zeta_n) - \sum_{j=1}^{k-1} \mu_j(\zeta_n) \nabla \mathcal{G}_j(\zeta_n).$$

Then, as  $n \to +\infty$ 

$$\mu_0(\zeta_n) \to \lambda_k \qquad \qquad \mu_j(\zeta_n) \to 0 \qquad (j = 1, \dots, k-1)$$

Moreover, if  $\xi_n = (\zeta_n)_{tr} \rightharpoonup 0$  in  $H^{1/2}$  then

$$\left(\xi_n, V\xi_n\right)_{L^2} \to 0.$$

*Proof.* We have that  $\zeta_n \in X_k$  is such that

$$\mathcal{I}(\zeta_n) \to \lambda_k \text{ and } \|\nabla_{X_k} \mathcal{I}(\zeta_n)\| \to 0.$$

Then  $\zeta_n$  is bounded in  $H^1$  and from (2.9) we have, as  $n \to +\infty$ 

$$\mu_0(\zeta_n) = \mathcal{I}(\zeta_n) - \frac{1}{2} (\nabla_{X_k} \mathcal{I}(\zeta_n), \zeta_n)_{H^1} \to \lambda_k.$$

Remark that, for all  $j \in \{1, \ldots, k-1\}$ 

$$0 = \nabla_{X_j} \mathcal{I}(\phi_j) = \nabla \mathcal{I}(\phi_j) - \mu_0(\phi_j) \nabla \mathcal{G}(\phi_j) - \sum_{i=1}^{j-1} \mu_i(\phi_j) \nabla \mathcal{G}_i(\phi_j).$$

and hence, for all  $\zeta_n \in X_k$  and  $j \in \{1, \ldots, k-1\}$  we have that

$$d\mathcal{I}(\zeta_{n})[\phi_{j}] = d\mathcal{I}(\phi_{j})[\zeta_{n}] = (\nabla \mathcal{I}(\phi_{j}), \zeta_{n})_{H^{1}}$$
$$= \mu_{0}(\phi_{j})(\nabla \mathcal{G}(\phi_{j}), \zeta_{n})_{H^{1}} + \sum_{i=1}^{j-1} \mu_{i}(\phi_{j})(\nabla \mathcal{G}_{i}(\phi_{j}), \zeta_{n})_{H^{1}} = 0$$

From this we conclude, using (2.10), that

$$\mu_j(\zeta_n) = d\mathcal{I}(\zeta_n)[\phi_j] - \left(\nabla_{X_k}\mathcal{I}(\zeta_n), \phi_j\right)_{H^1} = -\left(\nabla_{X_k}\mathcal{I}(\zeta_n), \phi_j\right)_{H^1} \to 0$$

for j = 1, ..., k - 1.

We then proceed as in the proof of Lemma 5. Since  $\zeta_n$  is a constrained Palais Smale sequence, we have

$$o_{n}(1) = C \|\nabla_{x_{k}} \mathcal{I}(\zeta_{n})\|_{H^{1}} \|\zeta_{n}\|_{H^{1}} \ge |(\nabla_{x_{k}} \mathcal{I}(\zeta_{n}), \chi_{R}^{2} \zeta_{n})|_{H^{1}}$$
  

$$\ge |d\mathcal{I}(\zeta_{n})[\chi_{R}^{2} \zeta_{n}]| - |\mu_{0}(\zeta_{n})(\nabla \mathcal{G}(\zeta_{n}), \chi_{R}^{2} \zeta_{n})_{H^{1}}|$$
  

$$- |\sum_{j=1}^{k-1} \mu_{j}(\zeta_{n})(\nabla \mathcal{G}_{j}(\zeta_{n}), \chi_{R}^{2} \zeta_{n})_{H^{1}}|$$
  

$$\ge 2\mathcal{I}(\chi_{R} \zeta_{n}) - 2c^{2} \|\zeta_{n} \nabla \chi_{R}\|_{L^{2}}^{2} - 2|\mu_{0}(\zeta_{n})||\chi_{R} \xi_{n}|_{L^{2}}^{2} - \sum_{j=1}^{k-1} |\mu_{j}(\zeta_{n})||\chi_{R} \xi_{n}|_{L^{2}}$$

where

$$\|\zeta_n \nabla \chi_{\scriptscriptstyle R}\|_{_{L^2}}^2 \leq C \sup_{y \in \mathbb{R}^3} |\nabla \chi_{\scriptscriptstyle R}|^2 \leq \frac{C}{R^2}$$

and by Sobolev compact embedding, for any given R > 0,

 $|\chi_{\scriptscriptstyle R}\xi_n|_{\scriptscriptstyle L^2} \to 0 \quad \text{as} \quad n \to +\infty.$ 

Moreover,  $|\mu_j(\zeta_n)| \leq C$  for  $j = 0, \ldots, k - 1$ .

Now, since  $\mathcal{I}$  is coercive, exactly as in Lemma 5 we may conclude

$$\|\chi_{R}\zeta_{n}\|_{H^{1}}^{2} \leq \epsilon_{n} + \frac{C}{R}$$

and by (h2) and Lemma 2,

$$|(V\chi_R\xi_n, \chi_R\xi_n)_{L^2}| \le a \|\chi_R\zeta_n\|_{H^1}^2 \le \epsilon_n + \frac{C}{R}$$

for  $\epsilon_n \to 0$  as  $n \to +\infty$ , R arbitrary large, and the Lemma follows.

We are now ready to prove the following Proposition for the existence of a minimizer for  $(\mathcal{P}_k)$ .

**Proposition 3.** Let  $\zeta_n \in X_k$  be a minimizing Palais Smale sequence for  $(\mathcal{P}_k)$ .

Then  $\zeta_n \rightharpoonup \phi_k$  in  $H^1$  and  $|(\phi_k)_{tr}|_{L^2}^{-1} \phi_k \in X_k$  is a minimizer for problem  $(\mathcal{P}_k)$ , and a weak solution of the Neumann problem  $(\mathcal{E}_k)$ .

*Proof.* We proceed as in the proof of Lemma 2 to conclude that  $\zeta_n \rightharpoonup \phi_k \neq 0$ .

We clearly have that  $\mathcal{G}_j(\phi_k) = 0$  for  $j = 1, \ldots, k-1$ . We do not know if  $|\varphi_k|_{L^2} = 1$  (where  $\varphi_k = (\phi_k)_{tr}$ ).

By Lemma 7 we have that

$$\mu_0(\zeta_n) \to \lambda_k \qquad \mu_j(\zeta_n) \to 0 \qquad (j = 1, \dots, k-1)$$

then by weak convergence we then have that for all  $h \in H^1$ , as  $n \to +\infty$ 

$$\left(\nabla_{X_k} \mathcal{I}(\zeta_n), h\right)_{H^1} = d\mathcal{I}(\zeta_n)[h] - 2\mu_0(\zeta_n) \operatorname{Re}(\xi_n, h_{tr})_{L^2} - \sum_{j=1}^{k-1} \mu_j(\zeta_n)(\varphi_j, h_{tr})_{L^2} \rightarrow d\mathcal{I}(\phi_k)[h] - 2\lambda_k \operatorname{Re}(\varphi_k, h_{tr})_{L^2} = 0.$$

We deduce, taking  $h = \phi_k$ 

$$0 = d\mathcal{I}(\phi_k) [\phi_k] - 2\lambda_k |\varphi_k|_{L^2}^2 = 2\mathcal{I}(\phi_k) - 2\lambda_k |\varphi_k|_{L^2}^2$$

and we conclude that  $|\varphi_k|_{L^2}^{-1}\phi_k \in X_k$  is a minimizer for  $(\mathcal{P}_k)$ .

Remark 6. It follows from the above Theorem that

(2.11) 
$$\nabla_{X_k} \mathcal{I}(\phi_k) = \nabla \mathcal{I}(\phi_k) - \lambda_k \nabla \mathcal{G}(\phi_k) = 0.$$

To conclude the proof of Theorem 1 we prove that  $\{\lambda_k\}_{k\geq 1} \in \sigma_{disc}(H_0 + V)$  namely that  $\lambda_k$  has finite multiplicity.

Indeed suppose that there exists an eigenvalue  $\lambda_k$  with infinite multiplicity. Then there exists a corresponding sequence  $\{\varphi_n^{(k)}\}_{n\in\mathbb{N}} \subset H^{1/2}$  of eigenfunctions corresponding to the same eigenvalue  $\lambda_k$ . We will assume that  $|\varphi_n^{(k)}|_{L^2} = 1$  for all  $n \in \mathbb{N}$ . Letting

$$\phi_n^{(k)} = \mathcal{F}_y^{-1} \left[ e^{-x\sqrt{m^2 c^4 + c^2 |p|^2}} \mathcal{F}[\varphi_n^{(k)}] \right] \in X_k,$$

by Lemma 2 we have  $\mathcal{I}(\phi_n^{(k)}) = \lambda_k$  and  $\nabla_{X_k} \mathcal{I}(\phi_n^{(k)}) = 0$ . We deduce from this that  $\varphi_n^k$  is a bounded sequence in  $H^{1/2}$ , since by orthogonality  $\varphi_n^{(k)} \to 0$  in  $L^2$ , we have  $\varphi_n^{(k)} \to 0$  in  $H^{1/2}$ , therefore by Lemma 5 we get

$$(\varphi_n^{(k)}, V\varphi_n^{(k)})_{L^2} \to 0 \quad \text{as } n \to +\infty$$

and from this we get a contradiction, namely  $\lambda_k = \mathcal{I}(\phi_n^{(k)}) \ge mc^2$ .

Finally since eigenvalues can accumulate only on the essential spectrum, we may conclude that

$$0 < \lambda_1 \le \dots \le \lambda_{k-1} \le \lambda_k \to \inf\{\sigma_{ess}(v)\} = mc^2 \text{ for } k \to +\infty$$

## 3. Proof of Theorem 2

Take  $\phi_k$  (and  $\varphi_k = (\phi_k)_{tr}$ ) and  $\lambda_k$  as in Theorem 1, and take R > 0 and T > 0, we set  $\chi_T(y) = \xi_R(y)g_T(y)$  where  $\xi_R(y) = \min\{(|y| - R)_+, 1\}$  and  $g_T(y) = \min\{e^{\frac{\beta}{c}|y|}, T\}$ , we introduce also the sets  $\mathcal{C}_R = \{(x, y) \in \mathbb{R}^4_+ \mid R < |y| < R + 1\}$  and  $\mathcal{D}_T = \{(x, y) \in \mathbb{R}^4_+ \mid e^{\frac{\beta}{c}|y|} < T\}$  where  $\xi_R$  and  $g_T$  are respectively not constants.

From (2.11), (2.8) and (2.3) we have

$$\begin{split} 0 &= \frac{1}{2} d\mathcal{I}(\phi_{k}) [\chi_{T}^{2} \phi_{k}] - \lambda_{k} \operatorname{Re} \left(\varphi_{k}, (\chi_{T}^{2} \phi_{k})_{tr}\right)_{L^{2}} \\ &= \frac{1}{2} d\mathcal{I}(\chi_{T} \phi_{k}) [\chi_{T} \phi_{k}] - c^{2} ||\phi_{k} \nabla_{y} \chi_{T}||_{L^{2}}^{2} - \lambda_{k} |\chi_{T} \varphi_{k}|_{L^{2}}^{2} \\ &= \mathcal{I}(\chi_{T} \phi_{k}) - c^{2} ||\phi_{k} \nabla_{y} \chi_{T}||_{L^{2}}^{2} - \lambda_{k} |\chi_{T} \varphi_{k}|_{L^{2}}^{2} \\ &= \mathcal{I}(\chi_{T} \phi_{k}) - \lambda_{k} |\chi_{T} \varphi_{k}|_{L^{2}}^{2} \\ &- \beta^{2} \iint_{\mathcal{D}_{T}} |\chi_{T} \phi_{k}|^{2} - c^{2} \iint_{\mathcal{C}_{R}} |\nabla_{y} \xi_{R}|^{2} |g_{T} \phi_{k}|^{2} \\ &- 2c\beta \iint_{\mathcal{D}_{T} \cap \mathcal{C}_{R}} \frac{y}{|y|} \cdot (\nabla_{y} \xi_{R}) \xi_{R} |g_{T} \phi_{k}|^{2} \\ &\geq \iint_{\mathbb{R}^{4}_{+}} |\partial_{x}(\chi_{T} \phi_{k})|^{2} + c^{2} \iint_{\mathbb{R}^{4}_{+}} |\nabla_{y}(\chi_{T} \phi_{k})|^{2} + (m^{2}c^{4} - \beta^{2}) \iint_{\mathbb{R}^{4}_{+}} |\chi_{T} \phi_{k}|^{2} \\ &- \int_{\mathbb{R}^{3}} |V| |\chi_{T} \varphi_{k}|^{2} - \lambda_{k} \int_{\mathbb{R}^{3}} |\chi_{T} \varphi_{k}|^{2} \\ &- c^{2} \iint_{\mathcal{C}_{R}} |g_{T} \phi_{k}|^{2} - 2c\beta \iint_{\mathcal{D}_{T} \cap \mathcal{C}_{R}} |g_{T} \phi_{k}|^{2} \\ &\geq \left( \sqrt{m^{2}c^{4} - \beta^{2}} - \lambda_{k} - \sup_{|y| \geq R} |V(y)| \right) \iint_{\mathbb{R}^{3}} |\chi_{T} \varphi_{k}|^{2} - C(R) \end{split}$$

Then, given  $\beta < \sqrt{m^2 c^4 - \lambda_k^2}$  there exists R > 0 such that

$$\sqrt{m^2c^4 - \beta^2} - \lambda_k - \sup_{|y| \ge R} |V(y)| > 0$$

and hence

$$\int_{\mathbb{R}^3} |\chi_T \varphi_k|^2 \le C$$

with C independent on T. Using monotone convergence we can pass to the limit as  $T \to +\infty$ to get

$$\int_{|y|\ge R} |\mathrm{e}^{\frac{\beta}{c}|y|}\varphi_k|^2 \le C$$

Namely,  $e^{\frac{\beta}{c}|y|}\varphi_k \in L^2(\mathbb{R}^3 \setminus B_R).$ 

### 4. Proof of Theorem 3

We need the following preliminary results.

**Proposition 4.** Let  $\phi_k \in H^1(\mathbb{R}^4_+)$  (and  $\varphi_k = (\phi_k)_{tr}$ ) as in Theorem 1 and  $V \in L^3_{loc}(\mathcal{U})$ . Then  $\phi_k \in L^p(\mathbb{R}_+ \times \mathcal{V})$  and  $\varphi_k \in L^p(\mathcal{V})$  for any  $p \geq 2$  and  $\mathcal{V} \subset \subset \mathcal{U}$ .

*Proof.* Take  $\phi_k$  (and  $\varphi_k = (\phi_k)_{tr}$ ) and  $\lambda_k$  as in Theorem 1, let  $v = \operatorname{Re} \phi_k$ . Take  $r, \delta > 0$ and  $y_0 \in \mathcal{U}$  such that the set  $B_{r+\delta}(y_0) = \{ y \in \mathbb{R}^3 \mid |y - y_0| \le r + \delta \} \subset \mathcal{U}$ . For  $n \in \mathbb{N}$ , let  $\xi_n(y) \in [0,1]$  a cut off function radial, piecewise linear and such that  $\xi_n(y) = 1$  if  $|y| \le r + \delta(\frac{2}{3})^n$  and  $\xi_n(y) = 0$  if  $|y| \ge r + \delta(\frac{2}{3})^{n-1}$ .

Let T > 0, we set  $v_T = \min\{v_+, T\}, \xi_n^0(y) = \xi_n(y - y_0)$  and  $\chi_{n,T}(x, y) = \xi_n^0(y)v_T^{\beta_n}(x, y)$ where  $\beta_n = (\frac{3}{2})^n - 1$ . We introduce also the sets  $B_n^0 = \{y \in \mathbb{R}^3 \mid \xi_n^0(y) = 1\}, C_n^0 =$  $\left\{ y \in \mathbb{R}^3 \mid |\nabla_y \xi_n^0(y)| = \frac{2}{\delta} (\frac{3}{2})^n \right\} \text{ and } D_T = \left\{ (x, y) \in \mathbb{R}^4_+ \mid v_+(x, y) < T \right\}.$  We have  $B_r(y_0) \subset B_{n+1}^0 \subset B_n^0 \subset B_{r+\delta}(y_0) \text{ and } C_{n+1}^0 \subset B_n^0 \text{ for any } n \in \mathbb{N}.$ From

$$\begin{split} \iint_{\mathbb{R}^{4}_{+}} |\partial_{x}(\chi_{n,T}v)|^{2} &= \iint_{\mathbb{R}^{4}_{+}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\partial_{x}v|^{2} + \beta_{n}^{2} \iint_{D_{T}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\partial_{x}v|^{2} \\ &+ 2\beta_{n} \iint_{D_{T}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\partial_{x}v|^{2} \\ &\geq (1+\beta_{n})^{2} \iint_{D_{T}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\partial_{x}v|^{2} \end{split}$$

we obtain

(4.1) 
$$\beta_n^2 \iint_{D_T} (\xi_n^0)^2 v_T^{2\beta_n} |\partial_x v|^2 \le \frac{\beta_n^2}{(1+\beta_n)^2} \iint_{\mathbb{R}^4_+} |\partial_x (\chi_{n,T} v)|^2$$

while from

$$\begin{split} &\iint_{\mathbb{R}^{4}_{+}} |\nabla_{y}(\chi_{n,T}v)|^{2} = \iint_{\mathbb{R}^{4}_{+}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\nabla_{y}v|^{2} + \beta_{n}^{2} \iint_{D_{T}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\nabla_{y}v|^{2} \\ &\quad + \iint_{\mathbb{R}^{4}_{+}} |\nabla_{y}\xi_{n}^{0}|^{2} |v_{T}^{\beta_{n}}v|^{2} + 2\beta_{n} \iint_{D_{T}} \xi_{n}^{0} (\nabla_{y}\xi_{n}^{0}) \cdot (\nabla_{y}v_{T}) v_{T}^{2\beta_{n}-1} |v|^{2} \\ &\quad + 2 \iint_{\mathbb{R}^{4}_{+}} \xi_{n}^{0} v_{T}^{2\beta_{n}} v \nabla_{y}\xi_{n}^{0} \cdot \nabla_{y}v + 2\beta_{n} \iint_{D_{T}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\nabla_{y}v|^{2} \\ &= \iint_{\mathbb{R}^{4}_{+} \setminus D_{T}} v_{T}^{2\beta_{n}} (v \nabla_{y}\xi_{n}^{0} + \xi_{n}^{0} \nabla_{y}v)^{2} + (\beta_{n}+1)^{2} \iint_{D_{T}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\nabla_{y}v|^{2} \\ &\quad + 2(\beta_{n}+1) \iint_{D_{T}} \xi_{n}^{0} (\nabla_{y}\xi_{n}^{0}) \cdot (\nabla_{y}v_{T}) v_{T}^{2\beta_{n}-1} |v|^{2} \\ &\geq (\beta_{n}+1) \left(\beta_{n} \iint_{D_{T}} (\xi_{n}^{0})^{2} v_{T}^{2\beta_{n}} |\nabla_{y}v|^{2} + 2 \iint_{D_{T}} \xi_{n}^{0} (\nabla_{y}\xi_{n}^{0}) \cdot (\nabla_{y}v_{T}) v_{T}^{2\beta_{n}-1} |v|^{2} \right). \end{split}$$

we deduce

$$(4.2) \quad \beta_n^2 \iint_{D_T} (\xi_n^0)^2 v_T^{2\beta_n} |\nabla_y v|^2 + 2\beta_n \iint_{D_T} \xi_n^0 (\nabla_y \xi_n^0) \cdot (\nabla_y v_T) v_T^{2\beta_n - 1} |v|^2 \\ \leq \frac{\beta_n}{\beta_n + 1} \iint_{\mathbb{R}^4_+} |\nabla_y (\chi_{n,T} v)|^2$$

Computations similar to those at the beginning of section 3, (we recall that  $v={\rm Re}\,\phi_k),$  leads to

$$\begin{split} 0 &= \frac{1}{2} d\mathcal{I}(\phi_k) [\chi^2_{n,T} v] - \lambda_k \operatorname{Re} \left( \varphi_k, (\chi^2_{n,T} v)_{tr} \right)_{L^2} \\ &= \mathcal{I}(\chi_{n,T} v) - \| v \partial_x \chi_{n,T} \|_{L^2}^2 - c^2 \| v \nabla_y \chi_{n,T} \|_{L^2}^2 - \lambda_k |\chi_{n,T} v|_{L^2}^2 \\ &= \mathcal{I}(\chi_{n,T} v) - \lambda_k |\chi_{n,T} v|_{L^2}^2 - \beta_n^2 \iint_{\mathcal{D}_T} (\xi_n^0)^2 v_T^{2\beta_n} |\partial_x v_T|^2 \\ &- c^2 \beta_n^2 \iint_{\mathcal{D}_T} (\xi_n^0)^2 v_T^{2\beta_n} |\nabla_y v_T|^2 - c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_T^{\beta_n} v|^2 \\ &- 2c^2 \beta_n \iint_{D_T} \xi_n^0 (\nabla_y \xi_n^0) \cdot (\nabla_y v_T) v_T^{2\beta_n - 1} v^2 \end{split}$$

Therefore we get, using (4.1) and (4.2)

$$\begin{split} 0 &= \iint_{\mathbb{R}^4_+} |\partial_x(\chi_{n,T}v)|^2 + c^2 \iint_{\mathbb{R}^4_+} |\nabla_y(\chi_{n,T}v)|^2 + m^2 c^4 \iint_{\mathbb{R}^4_+} |\chi_{n,T}v|^2 \\ &\quad - \beta_n^2 \iint_{D_T} (\xi_n^0)^2 v_T^{2\beta_n} |\partial_x v_T|^2 - c^2 \beta_n^2 \iint_{D_T} (\xi_n^0)^2 v_T^{2\beta_n} |\nabla_y v_T|^2 \\ &\quad - c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_T^{\beta_n}v|^2 - 2c^2 \beta_n \iint_{D_T} \xi_n^0 (\nabla_y \xi_n^0) \cdot (\nabla_y v_T) v_T^{2\beta_n - 1} v^2 \\ &\quad + \int_{\mathbb{R}^3} V|(\chi_{n,T}v)_{tr}|^2 - \lambda_k \int_{\mathbb{R}^3} |(\chi_{n,T}v)_{tr}|^2 \\ &\geq \left(1 - \frac{\beta_n}{\beta_n + 1}\right) \|\chi_{n,T}v\|_{H^1}^2 - c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_T^{\beta_n}v|^2 \\ &\quad - \int_{\mathbb{R}^3} |V||(\chi_{n,T}v)_{tr}|^2 - \lambda_k \int_{\mathbb{R}^3} |(\chi_{n,T}v)_{tr}|^2 \end{split}$$

Namely,

$$\frac{1}{\beta_n+1} \|\xi_n^0 v_T^{\beta_n} v\|_{H^1}^2 \\
\leq c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_T^{\beta_n} v|^2 + \int_{\mathbb{R}^3} |V| |\xi_n^0 (v_T^{\beta_n} v)_{tr}|^2 + \lambda_k \int_{\mathbb{R}^3} |\xi_n^0 (v_T^{\beta_n} v)_{tr}|^2$$

Using Fatou's Lemma and monotone convergence, we can pass to the limit as  $T \to +\infty$ to get

$$(4.3) \quad \frac{1}{\alpha_n} \|\xi_n^0 v_+^{\alpha_n}\|_{H^1}^2 \\ \leq c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_+^{\alpha_n}|^2 + \int_{\mathbb{R}^3} |V| |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 + \lambda_k \int_{\mathbb{R}^3} |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 \\ \leq c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_+^{\alpha_n}|^2 + \int_{\mathbb{R}^3} |V| |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 + \lambda_k \int_{\mathbb{R}^3} |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 \\ \leq c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_+^{\alpha_n}|^2 + \int_{\mathbb{R}^3} |V| |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 + \lambda_k \int_{\mathbb{R}^3} |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 \\ \leq c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_+^0|^2 + \int_{\mathbb{R}^3_+} |V| |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 \\ \leq c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_+^0|^2 + \int_{\mathbb{R}^3_+} |V| |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 + \lambda_k \int_{\mathbb{R}^3_+} |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 \\ \leq c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n^0|^2 |v_+^0 (v_+)_{tr}^{\alpha_n}|^2 + \int_{\mathbb{R}^3_+} |V| |\xi_n^0 (v_+)_{tr}^{\alpha_n}|^2 \\ \leq c^2 \iint_{\mathbb{R}^4_+} |v_+^0 (v_+)_{tr}^{\alpha_n}|^2 + \int_{\mathbb{R}^3_+} |v_+|^2 |v_+|^2 + \int_{\mathbb{R}^3_+} |v_+|^2$$

where  $\alpha_n = \beta_n + 1 = (3/2)^n$ For any M > 0, let  $A_1 = \{|V| \le M\} \cap B_{r+\delta}(y_0), A_2 = \{|V| > M\} \cap B_{r+\delta}(y_0)$ , then, since  $V \in L^3_{loc}(\mathcal{U})$ , we have

$$\begin{split} \int_{\mathbb{R}^3} |V| |\xi_n^0(v_+)_{tr}^{\alpha_n}|^2 &\leq \int_{A_1} |V| |\xi_n^0(v_+)_{tr}^{\alpha_n}|^2 + \int_{A_2} |V| |\xi_n^0(v_+)_{tr}^{\alpha_n}|^2 \\ &\leq M \int_{A_1} |\xi_n^0(v_+)_{tr}^{\alpha_n}|^2 + \left(\int_{A_2} |V|^3\right)^{1/3} \left(\int_{A_2} |\xi_n^0(v_+)_{tr}^{\alpha_n}|^3\right)^{2/3} \\ &\leq M |\xi_n^0(v_+)_{tr}^{\alpha_n}|^2 + \epsilon(M) |\xi_n^0(v_+)_{tr}^{\alpha_n}|^3 \end{split}$$

now we take  $C \ge \max\{(\frac{2}{\delta}c)^2, M + \lambda_k\}$  and we get

$$\begin{aligned} \|\xi_{n}^{0}v_{+}^{\alpha_{n}}\|_{H^{1}}^{2} \leq C \left( \alpha_{n}^{3} \iint_{\mathbb{R}_{+} \times C_{n}^{0}} |v_{+}^{\alpha_{n}}|^{2} + \alpha_{n} |\xi_{n}^{0}(v_{+})_{tr}^{\alpha_{n}}|_{2}^{2} \right) \\ + \alpha_{n} \epsilon(M) |\xi_{n}^{0}(v_{+})_{tr}^{\alpha_{n}}|_{3}^{2} \end{aligned}$$

Taking M sufficiently large, that is  $\epsilon(M)$  sufficiently small, by Sobolev inequality we have

$$\|(v_{+}^{\alpha_{n}})_{tr}\|_{L^{2^{\sharp}}(B_{n}^{0})}^{2} + \|v_{+}^{\alpha_{n}}\|_{L^{2^{*}}(\mathbb{R}_{+}\times B_{n}^{0})}^{2} \leq K_{n} \left(\|v_{+}^{\alpha_{n}}\|_{L^{2}(\mathbb{R}_{+}\times C_{n}^{0})}^{2} + |\xi_{n}^{0}(v_{+}^{\alpha_{n}})_{tr}|_{2}^{2}\right)$$

where  $2^{\sharp} = 2N/(N-1) = 3$  (here N = 3) and  $2^* = 2N/(N-2) = 4$  (here N = 4) are the critical Sobolev exponent for the embedding of  $H^{1/2}(\mathbb{R}^3)$  in  $L^p(\mathbb{R}^3)$  and for the embedding of  $H^1(\mathbb{R}^4)$  in  $L^p(\mathbb{R}^4)$  and the constant  $K_n$  depend on  $n \in \mathbb{N}$ .

Finally, since  $C_n^0 \subset B_{n-1}^0$  we may conclude

(4.4) 
$$\begin{cases} |(v_{+})_{tr}^{\alpha_{n}}|^{2}_{L^{2^{\sharp}}(B_{n}^{0})} \leq K_{n} \left( \|v_{+}^{\alpha_{n}}\|^{2}_{L^{2}(\mathbb{R}_{+}\times B_{n-1}^{0})} + |(v_{+})_{tr}^{\alpha_{n}}|^{2}_{L^{2}(B_{n-1}^{0})} \right) \\ \|v_{+}^{\alpha_{n}}\|^{2}_{L^{2^{*}}(\mathbb{R}_{+}\times B_{n}^{0})} \leq K_{n} \left( \|v_{+}^{\alpha_{n}}\|^{2}_{L^{2}(\mathbb{R}_{+}\times B_{n-1}^{0})} + |(v_{+})_{tr}^{\alpha_{n}}|^{2}_{L^{2}(B_{n-1}^{0})} \right). \end{cases}$$

Then a bootstrap argument can start: since  $v_+ \in H^1(\mathbb{R}^4_+)$  we have  $v_+ \in L^p(\mathbb{R}^4_+)$  for  $p \in [2,4]$  and  $(v_+)_{tr} \in L^q(\mathbb{R}^3)$  for  $q \in [2,3]$ , hence we can apply (4.4) with n = 1 to deduce that  $(v_+)_{tr} \in L^{2^{\sharp \alpha_1}}(B_1^0) = L^{3(3/2)}(B_1^0)$  and  $v_+ \in L^{2^*\alpha_1}(\mathbb{R}_+ \times B_1^0) = L^6(\mathbb{R}_+ \times B_1^0)$ . Since  $2\alpha_n = 2^{\sharp \alpha_{n-1}} < 2^*\alpha_{n-1}$  we can then apply again (4.4) and, after n iterations, we deduce that  $(v_+)_{tr} \in L^{3(3/2)^n}(B_n^0)$ ,  $v_+ \in L^{4(3/2)^n}(\mathbb{R}_+ \times B_n^0)$ . Hence we may conclude that  $(v_+)_{tr} \in L^p(B_r(y_0))$  and  $v_+ \in L^p(\mathbb{R}_+ \times B_r(y_0))$  for all  $p \in [2, +\infty)$ .

The same is clearly true for  $v_{-}$  and hence for  $v = \operatorname{Re} \phi_k$ . Analogously we can argue for  $\operatorname{Im} \phi_k$  and we get the result for  $\varphi_k = (\phi_k)_{tr}$ .

**Proposition 5.** Let  $\phi_k \in H^1(\mathbb{R}^4_+)$  (and  $\varphi_k = (\phi_k)_{tr}$ ) as in Theorem 1. Then given any  $R > R_0$  (with  $R_0$  given in (h1)) we have  $\phi_k \in L^p(\mathbb{R}_+ \times (\mathbb{R}^3 \setminus B_R))$  and  $\varphi_k \in L^p(\mathbb{R}^3 \setminus B_R)$  for any  $p \in [2, \infty]$ .

*Proof.* By (**h1**) we have  $V \in L^{\infty}(\mathbb{R}^3 \setminus B_{R_0})$  for some  $R_0 > 0$ . Take  $\phi_k$  (and  $\varphi_k = (\phi_k)_{tr}$ ) and  $\lambda_k$  as in Theorem 1, let  $v = \operatorname{Re} \phi_k$ .

Take any  $\delta > 0$  and for  $n \in \mathbb{N}$  let  $\xi_n(y) \in [0, 1]$  be a cut off function, radial, piecewise linear and such that  $\xi_n(y) = 0$  if  $|y| \le R_0 + \delta \sum_{k=0}^{n-1} {\binom{2}{3}}^k$  and  $\xi_n(y) = 1$  if  $|y| \ge R_0 + \delta \sum_{k=0}^n {\binom{2}{3}}^k$ .

Let T > 0, we set  $v_T = \min\{v_+, T\}$  and  $\chi_{n,T}(x, y) = \xi_n(y)v_T^{\beta_n}(x, y)$  where  $\beta_n = (\frac{3}{2})^n - 1$ . We introduce also the sets  $F_n = \{y \in \mathbb{R}^3 \mid \xi_n(y) = 1\}, C_n = \{y \in \mathbb{R}^3 \mid |\nabla_y \xi_n^0(y)| = \frac{2}{\delta} (\frac{3}{2})^n\}$ and  $D_T = \{(x, y) \in \mathbb{R}^4_+ \mid v_+(x, y) < T\}$ . We have  $\mathbb{R}^3 \setminus B_{R_0+\delta} \subset F_{n+1} \subset F_n \subset \mathbb{R}^3 \setminus B_{R_0}$ and  $C_{n+1} \subset F_n$  for any  $n \in \mathbb{N}$ .

Now we can repeat the estimates in the proof of Proposition 4 to deduce that also in this case (4.3) holds, namely

$$\frac{1}{\alpha_n} \|\xi_n v_+^{\alpha_n}\|_{H^1}^2 \le c^2 \iint_{\mathbb{R}^4_+} |\nabla_y \xi_n|^2 |v_+^{\alpha_n}|^2 + \int_{\mathbb{R}^3} |V| |\xi_n (v_+^{\alpha_n})_{tr}|^2 + \lambda_k \int_{\mathbb{R}^3} |\xi_n (v_+^{\alpha_n})$$

where also here  $\alpha_n = \beta_n + 1 = (3/2)^n$ .

Then taking a positive constant  $C \ge \max\{(\frac{2}{\delta}c)^2, (\sup_{\mathbb{R}^3 \setminus B_{R_0}} |V| + \lambda_k)\}$  we get

$$\|\xi_n v_+^{\alpha_n}\|_{H^1}^2 \le C\left(\alpha_n^3 \iint_{\mathbb{R}_+ \times C_n} |v_+^{\alpha_n}|^2 + \alpha_n \int_{\mathbb{R}^3} |\xi_n (v_+^{\alpha_n})_{tr}|^2\right)$$

and again by Sobolev inequality and recalling that  $C_n \subset F_{n-1}$ 

$$|(v_{+})_{tr}^{\alpha_{n}}|_{L^{2^{\sharp}}(F_{n})}^{2} + ||v_{+}^{\alpha_{n}}||_{L^{2^{*}}(\mathbb{R}_{+}\times F_{n})}^{2} \leq C\left(\alpha_{n}^{3}||v_{+}^{\alpha_{n}}||_{L^{2}(\mathbb{R}_{+}\times F_{n-1})}^{2} + \alpha_{n}|(v_{+})_{tr}^{\alpha_{n}}|_{L^{2}(F_{n-1})}^{2}\right)$$

Finally, we may conclude

(4.5) 
$$\begin{cases} |(v_{+})_{tr}^{\alpha_{n}}|_{L^{2^{\sharp}}(F_{n})}^{2} \leq C\left(\alpha_{n}^{3}\|v_{+}^{\alpha_{n}}\|_{L^{2}(\mathbb{R}_{+}\times F_{n-1})}^{2} + \alpha_{n}|(v_{+})_{tr}^{\alpha_{n}}|_{L^{2}(F_{n-1})}^{2}\right) \\ \|v_{+}^{\alpha_{n}}\|_{L^{2^{*}}(\mathbb{R}_{+}\times F_{n})}^{2} \leq C\left(\alpha_{n}^{3}\|v_{+}^{\alpha_{n}}\|_{L^{2}(\mathbb{R}_{+}\times F_{n-1})}^{2} + \alpha_{n}|(v_{+})_{tr}^{\alpha_{n}}|_{L^{2}(F_{n-1})}^{2}\right). \end{cases}$$

Then, exactly as in the proof of Proposition 4, a bootstrap argument can start and after n iterations, we deduce that  $(v_+)_{tr} \in L^{3(3/2)^n}(F_n)$ ,  $v_+ \in L^{4(3/2)^n}(\mathbb{R}_+ \times F_n)$ . Hence we may conclude that  $(v_+)_{tr} \in L^p(\mathbb{R}^3 \setminus B_{R_0+\delta})$  and  $v_+ \in L^p(\mathbb{R}_+ \times (\mathbb{R}^3 \setminus B_{R_0+\delta}))$  for all  $p \in [2, \infty)$ .

To prove that actually  $(v_+)_{tr} \in L^{\infty}(\mathbb{R}^3 \setminus B_{R_0+\delta})$  and  $v_+ \in L^{\infty}(\mathbb{R}_+ \times (\mathbb{R}^3 \setminus B_{R_0+\delta}))$  we can argue as follows. In view of (4.5) we have

$$\begin{aligned} |(v_{+})_{tr}|_{L^{2^{\sharp}\alpha_{n}}(F_{n})}^{2\alpha_{n}} \leq & C\left(\alpha_{n}^{3}\|v_{+}\|_{L^{2\alpha_{n}}(\mathbb{R}_{+}\times F_{n-1})}^{2\alpha_{n}} + \alpha_{n}|(v_{+})_{tr}|_{L^{2\alpha_{n}}(F_{n-1})}^{2\alpha_{n}}\right) \\ \leq & M_{0}^{2} e^{2\sqrt{\alpha_{n}}} \left(\max\{\|v_{+}\|_{L^{2\alpha_{n}}(\mathbb{R}_{+}\times F_{n-1})}, |(v_{+})_{tr}|_{L^{2\alpha_{n}}(F_{n-1})}\}\right)^{2\alpha_{n}} \end{aligned}$$

Moreover, since

$$\|v_{+}^{\alpha_{n}}\|_{L^{2^{\sharp}}} \leq \|v_{+}^{\alpha_{n}}\|_{L^{2}}^{1/2} \|v_{+}^{\alpha_{n}}\|_{L^{2^{\ast}}}^{1/2}$$

and  $F_n \subset F_{n-1}$  we have

$$\|v_{+}\|_{L^{2^{\sharp}\alpha_{n}}(\mathbb{R}_{+}\times F_{n})}^{2\alpha_{n}} \leq \|v_{+}\|_{L^{2\alpha_{n}}}^{\alpha_{n}} \|v_{+}\|_{L^{2^{\ast}\alpha_{n}}}^{\alpha_{n}} \leq \frac{1}{2} \|v_{+}\|_{L^{2\alpha_{n}}}^{2\alpha_{n}} + \frac{1}{2} \|v_{+}\|_{L^{2^{\ast}\alpha_{n}}}^{2\alpha_{n}} \\ \leq M_{0}^{2} e^{2\sqrt{\alpha_{n}}} \left( \max\{\|v_{+}\|_{L^{2\alpha_{n}}(\mathbb{R}_{+}\times F_{n-1})}, |(v_{+})_{tr}|_{L^{2\alpha_{n}}(F_{n-1})} \} \right)^{2\alpha_{n}}$$

where the positive constant  $M_0 > 1$  is independent of n. Hence, recalling also that  $2^{\sharp} \alpha_n = 2\alpha_{n+1}$ , we get

$$\begin{cases} |(v_{+})_{tr}|_{L^{2\alpha_{n+1}}(F_{n})} &\leq M_{0}^{\frac{1}{\alpha_{n}}} e^{\frac{1}{\sqrt{\alpha_{n}}}} \max\{\|v_{+}\|_{L^{2\alpha_{n}}(\mathbb{R}_{+}\times F_{n-1})}, |(v_{+})_{tr}|_{L^{2\alpha_{n}}(F_{n-1})}\} \\ \|v_{+}\|_{L^{2\alpha_{n+1}}(\mathbb{R}_{+}\times F_{n})} &\leq M_{0}^{\frac{1}{\alpha_{n}}} e^{\frac{1}{\sqrt{\alpha_{n}}}} \max\{\|v_{+}\|_{L^{2\alpha_{n}}(\mathbb{R}_{+}\times F_{n-1})}, |(v_{+})_{tr}|_{L^{2\alpha_{n}}(F_{n-1})}\} \end{cases}$$

We set  $A_n = \max\{\|v_+\|_{L^{2\alpha_n}(\mathbb{R}_+ \times F_{n-1})}, |(v_+)_{tr}|_{L^{2\alpha_n}(F_{n-1})}\}$  then we have

$$A_{n+1} \le M_0^{\frac{1}{\alpha_n}} e^{\frac{1}{\sqrt{\alpha_n}}} A_n \le M_0^{\sum_{i=0}^n \frac{1}{\alpha_i}} e^{\sum_{i=0}^n \frac{1}{\sqrt{\alpha_i}}} A_0.$$

Since

$$\sum_{i=0}^{+\infty} \frac{1}{\sqrt{\alpha_i}} < +\infty$$

then there exists a constant K independent on p such that  $|(v_+)_{tr}|_{L^p(\mathbb{R}^3 \setminus B_{R_0+\delta})} < K$  and  $||v_+||_{L^p(\mathbb{R}_+ \times (\mathbb{R}^3 \setminus B_{R_0+\delta}))} < K$ , for any  $p \ge 2$  and we deduce that  $(v_+)_{tr} \in L^{\infty}(\mathbb{R}^3 \setminus B_{R_0+\delta})$  and  $v_+ \in L^{\infty}(\mathbb{R}_+ \times (\mathbb{R}^3 \setminus B_{R_0+\delta}))$ .

The same is clearly true for  $v_{-}$  and hence for  $v = \operatorname{Re} \phi_k$ . Analogously we can argue for  $\operatorname{Im} \phi_k$  and we get the result for  $\varphi_k = (\phi_k)_{tr}$ .

Now we finally conclude the proof of Theorem 3 as follow

(i): Recalling that  $\phi_k \in H^1(\mathbb{R}^4_+, \mathbb{C})$  is a weak solution of the Neumann problem

$$\begin{cases} -\partial_x^2 \phi_k - c^2 \Delta_y \phi_k + m^2 c^4 \phi_k = 0 & \text{ in } \mathbb{R}^4_+ \\ \frac{\partial \phi_k}{\partial \nu} + V \varphi_k = \lambda_k \varphi_k & \text{ on } \partial \mathbb{R}^4_+ = \mathbb{R}^3. \end{cases}$$

then following [2] we introduce

$$\psi_k(x,y) = \int_0^x \phi_k(t,y) \, dt$$

we clearly have that  $\psi_k \in H^1((0, r) \times \mathbb{R}^3, \mathbb{C})$  for any r > 0 and we have (see [5, Proposition 3.9] for the details) that  $\psi_k$  is a weak solution of the following Dirichlet problem

$$\begin{cases} -\partial_x^2 \psi_k - c^2 \Delta_y \psi_k + m^2 c^4 \psi_k = f(x, y) & \text{in } \mathbb{R}^4_+ \\ \psi_k = 0 & \text{on } \partial \mathbb{R}^4_+ = \mathbb{R}^3 \end{cases}$$

where  $f(x, y) = (\lambda_k - V(y))\varphi_k(y)$ .

Now let us define

$$(\psi_k)_{odd}(x,y) = \begin{cases} \psi_k(x,y) & x \ge 0\\ -\psi_k(-x,y) & x < 0 \end{cases} \quad \text{and} \quad f_{odd}(x,y) = \begin{cases} f(x,y) & x \ge 0\\ -f(x,y) & x < 0 \end{cases}$$

It is easy to check that  $(\psi_k)_{odd} \in H^1((-r,r) \times \mathbb{R}^3, \mathbb{C})$  is a weak solution of the (linear) second order elliptic problem

$$-\partial_x^2 u - c^2 \Delta_y u + m^2 c^4 u = f_{odd} \quad \text{in } \mathbb{R}^4.$$

Since by Proposition 5  $f_{odd} \in L^q((-r,r) \times (\mathbb{R}^3 \setminus B_R))$  for any  $q \in [2,\infty]$ , r > 0 and  $R > R_0$ we deduce by standard elliptic regularity that  $(\psi_k)_{odd} \in W^{2,q}((-r,r) \times (\mathbb{R}^3 \setminus B_R))$  and hence in particular  $\phi_k = \partial_x \psi_k \in W^{1,q}((0,r) \times (\mathbb{R}^3 \setminus B_R))$ .

(ii) : By Sobolev's embedding  $\psi_k \in C^{1,\alpha}([0,+\infty) \times (\mathbb{R}^3 \setminus B_R))$  for all  $\alpha \in [0,1]$ . Namely, we get that  $\phi_k = \partial_x \psi_k \in C^{0,\alpha}([0,+\infty) \times (\mathbb{R}^3 \setminus B_R))$  and  $\varphi_k = \phi_k(0, \cdot) \in C^{0,\alpha}(\mathbb{R}^3 \setminus B_R)$  for any  $\alpha \in [0,1]$  and  $R > R_0$ .

(iii): Since by Proposition 4  $f_{odd} \in L^q((-r,r) \times \mathcal{V})$  for any  $q \in [2,\infty)$ , r > 0 and  $\mathcal{V} \subset \subset \mathcal{U}$ we deduce by standard elliptic regularity that  $(\psi_k)_{odd} \in W^{2,q}((-r,r) \times \mathcal{V})$  hence in particular  $\phi_k = \partial_x \psi_k \in W^{1,q}((0,r) \times \mathcal{V})$ . Then by the trace Theorem we get  $\varphi_k \in W^{1-\frac{1}{q},q}(\mathcal{V})$  for any  $q \in [2,\infty)$  and  $\mathcal{V} \subset \subset \mathcal{U}$  and by Sobolev embedding  $\varphi_k \in C^{0,\alpha}(\mathcal{V})$  for any  $\alpha \in [0,1)$ .

### References

- S. Agmon, Lectures on exponential decay of solutions of second-order elliptic equations: bounds on eigenfunctions of N-body Schrödinger operators, Mathematical Notes, vol. 29, Princeton University Press, Princeton, NJ; University of Tokyo Press, Tokyo, 1982.
- [2] X. Cabré and J. Solà-Morales, Layer solutions in a half-space for boundary reactions, Comm. Pure Appl. Math. 58 (2005), no. 12, 1678–1732.
- [3] L. Caffarelli and L. Silvestre, An extension problem related to the fractional Laplacian, Comm. Partial Differential Equations 32 (2007), no. 7-9, 1245–1260.
- [4] R. Carmona, W. C. Masters, and B. Simon, Relativistic Schrödinger operators: asymptotic behavior of the eigenfunctions, J. Funct. Anal. 91 (1990), no. 1, 117–142.
- [5] V. Coti Zelati and M. Nolasco, Existence of ground states for nonlinear, pseudorelativistic Schrödinger equations, Rend. Lincei Mat. Appl. 22 (2011), 51–72.

- [6] \_\_\_\_\_, Ground states for pseudo-relativistic Hartree equations of critical type, Rev. Mat. Iberoam. 29 (2013), no. 4, 1421–1436.
- [7] \_\_\_\_\_, A variational approach to the Brown-Ravenhall operator for the relativistic one-electron atoms, preprint, 2014, arXiv:1103.2649.
- [8] J. Dolbeault, M. J. Esteban, and E. Séré, Variational characterization for eigenvalues of Dirac operators, Calc. Var. Partial Differential Equations 10 (2000), no. 4, 321–347.
- M. J. Esteban, M. Lewin, and E. Séré, Variational methods in relativistic quantum mechanics, Bull. Amer. Math. Soc. (N.S.) 45 (2008), no. 4, 535–593.
- [10] I. W. Herbst, Spectral theory of the operator  $(p^2 + m^2)^{1/2} Ze^2/r$ , Comm. Math. Phys. 53 (1977), no. 3, 285–294.
- [11] V. Kovalenko, M. Perelmuter, and Y. Semenov, Schrödinger operators with potentials, J.Math.Phys. 22 (1981), no. 5, 1033–1044.
- [12] E. Lenzmann, Well-posedness for semi-relativistic Hartree equations of critical type, Math. Phys. Anal. Geom. 10 (2007), no. 1, 43–64.
- [13] E. H. Lieb and M. Loss, Analysis, Graduate Studies in Mathematics, no. 14, American Mathematical Society, 1997.
- [14] F. Nardini, Exponential decay for the eigenfunctions of the two-body relativistic Hamiltonian, J. Analyse Math. 47 (1986), 87–109.
- [15] A. J. O'Connor, Exponential decay of bound state wave functions, Comm. Math. Phys. 32 (1973), 319–340.
- [16] B. Simon, Fifty years of eigenvalue perturbation theory, AMS-MAA Joint Lecture Series, American Mathematical Society, Providence, RI, 1990, A joint AMS-MAA lecture presented in Louisville, Kentucky, January 1990.

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